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## Supplementary Information for: Ab-initio calculations of temperature dependent electronic structure of inorganic halide perovskite materials

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Fig. S1. The dependence of band energy renormalization in the OTMS approach on the number of q-points and energy level broadening parameter  $\delta$  for CsPbCl<sub>3</sub> at T = 400 K.



Fig. S2. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the R point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S3. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the R point.



Fig. S4. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the R point in the case of  $CsPbCl_3$  material at T = 400 K.



Fig. S5. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the R point.



Fig. S6. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the R point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S7. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the R point.



Fig. S8. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbCl<sub>3</sub> at the R point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 320$  K and band energy from PBE0 calculations, respectively.



Fig. S9. The dependence of band energy renormalization in the OTMS approach on the number of q-points and energy level broadening parameter  $\delta$  for CsPbBr<sub>3</sub> at T = 400 K.



Fig. S10. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the R point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S11. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the R point.



Fig. S12. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the R point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S13. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the R point.



Fig. S14. The dependence of band energy renormalization in the OTMS approach on the number of q-points and energy level broadening parameter  $\delta$  for CsPbI<sub>3</sub> at T = 400 K.



Fig. S15. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the R point in the case of CsPbI<sub>3</sub> material at T = 400 K.



Fig. S16. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the R point.



Fig. S17. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the R point in the case of CsPbI<sub>3</sub> material at T = 400 K.



Fig. S18. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the R point.



Fig. S19. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the R point in the case of  $CsPbI_3$  material at T = 400 K.



Fig. S20. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the R point.



Fig. S21. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbI<sub>3</sub> at the R point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 300$  K and band energy from PBE0 calculations, respectively.



Fig. S22. Contributions from Fan and Debye-Waller terms of phonons at different frequencies to VBM (left column) and CBM (right column) renormalization at T = 400 K for CsPbX<sub>3</sub> (X = Cl, Br or I, in rows from top to bottom) obtained using the OTMS approach. The value of each bin  $b(\omega_i)$  represents the contribution of all phonons with frequencies from the range  $(\omega_i - \Delta\omega/2, \omega_i + \Delta\omega/2)$  to band energy renormalization, so that  $\Delta E_{\mathbf{kn}} = \sum_i b(\omega_i)$ .



Fig. S23. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the  $\Gamma$  point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S24. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the  $\Gamma$  point.



Fig. S25. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the  $\Gamma$  point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S26. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the  $\Gamma$  point.



Fig. S27. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the  $\Gamma$  point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S28. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the  $\Gamma$  point.



Fig. S29. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbCl<sub>3</sub> at the  $\Gamma$  point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 320$  K and band energy from PBE0 calculations, respectively.



Fig. S30. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the  $\Gamma$  point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S31. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the  $\Gamma$  point.



Fig. S32. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the  $\Gamma$  point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S33. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the  $\Gamma$  point.



Fig. S34. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the  $\Gamma$  point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S35. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the  $\Gamma$  point.



Fig. S36. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbBr<sub>3</sub> at the  $\Gamma$  point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 403$  K and band energy from PBE0 calculations, respectively.



Fig. S37. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the  $\Gamma$  point in the case of CsPbI<sub>3</sub> material at T = 400 K.



Fig. S38. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the  $\Gamma$  point.



Fig. S39. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the  $\Gamma$  point in the case of CsPbI<sub>3</sub> material at T = 400 K.



Fig. S40. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the  $\Gamma$  point.



Fig. S41. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the  $\Gamma$  point in the case of CsPbI<sub>3</sub> material at T = 400 K.



Fig. S42. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the  $\Gamma$  point.



Fig. S43. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbI<sub>3</sub> at the  $\Gamma$  point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 300$  K and band energy from PBE0 calculations, respectively.



Fig. S44. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the X point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S45. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the X point.



Fig. S46. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the X point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S47. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the X point.



Fig. S48. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the X point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S49. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the X point.



Fig. S50. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbCl<sub>3</sub> at the X point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 320$  K and band energy from PBE0 calculations, respectively.



Fig. S51. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the X point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S52. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the X point.



Fig. S53. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the X point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S54. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the X point.



Fig. S55. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the X point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S56. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the X point.



Fig. S57. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbBr<sub>3</sub> at the X point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 403$  K and band energy from PBE0 calculations, respectively.



Fig. S58. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the X point in the case of CsPbI<sub>3</sub> material at T = 400 K.



Fig. S59. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the X point.



Fig. S60. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the X point in the case of CsPbI<sub>3</sub> material at T = 400 K.



Fig. S61. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the X point.



Fig. S62. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the X point in the case of  $CsPbI_3$  material at T = 400 K.



Fig. S63. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the X point.



Fig. S64. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbI<sub>3</sub> at the X point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 300$  K and band energy from PBE0 calculations, respectively.



Fig. S65. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the M point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S66. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the M point.



Fig. S67. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the M point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S68. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the M point.



Fig. S69. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the M point in the case of CsPbCl<sub>3</sub> material at T = 400 K.



Fig. S70. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbCl<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the M point.



Fig. S71. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbCl<sub>3</sub> at the M point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 320$  K and band energy from PBE0 calculations, respectively.



Fig. S72. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the M point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S73. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the M point.



Fig. S74. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the M point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S75. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the M point.



Fig. S76. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the M point in the case of CsPbBr<sub>3</sub> material at T = 400 K.



Fig. S77. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbBr<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the M point.



Fig. S78. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbBr<sub>3</sub> at the M point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 403$  K and band energy from PBE0 calculations, respectively.



Fig. S79. The frequency dependence of the self-energy and the spectral function for bands VBM4, VBM3, and VBM2 at the M point in the case of CsPbI<sub>3</sub> material at T = 400 K.



Fig. S80. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands VBM4, VBM3, and VBM2 at the M point.



Fig. S81. The frequency dependence of the self-energy and the spectral function for bands VBM1, VBM, CBM and CBM1 at the M point in the case of CsPbI<sub>3</sub> material at T = 400 K.



Fig. S82. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands VBM1, VBM, CBM and CBM1 at the M point.



Fig. S83. The frequency dependence of the self-energy and the spectral function for bands CBM2, CBM3, and CBM4 at the M point in the case of CsPbI<sub>3</sub> material at T = 400 K.



Fig. S84. The dependence of the real and imaginary part of self-energy at the renormalized energy on the size of the q-points grid. The results are presented for CsPbI<sub>3</sub> material at T = 400 K for bands CBM2, CBM3, and CBM4 at the M point.



Fig. S85. Temperature dependence of the band energy and the imaginary part of the self-energy for VBM and VBMx (CBM and CBMx) bands (where x = 1, 2, 3, 4) calculated using the SCP approach. The results are shown for CsPbI<sub>3</sub> at the M point. Vertical and horizontal dotted lines represent the temperature of the phase transition to cubic structure  $T_c = 300$  K and band energy from PBE0 calculations, respectively.



Fig. S86. Electronic band structure obtained from DFT using PBEsol functional for  $CsPbCl_3$  and  $CsPbBr_3$ , and PBE functional for  $CsPbI_3$ .



Fig. S87. Spectral function of CsPbCl<sub>3</sub> at T = 50 K and T = 100 K for VBM2 band at the X point. The inset shows zoomed figure around the peaks to show the shift of spectral function maximum with temperature.



Fig. S88. Spectral function of CsPbBr<sub>3</sub> at T = 600 K and T = 650 K for VBM1 band at the X point. The inset shows zoomed figure around the peaks to show the shift of spectral function maximum with temperature.



Fig. S89. Spectral function of CsPbI<sub>3</sub> at T = 550 K and T = 600 K for VBM1 band at the X point. The inset shows zoomed figure around the peaks to show the shift of spectral function maximum with temperature.