

# Bose-Einstein Condensation in Random Potentials

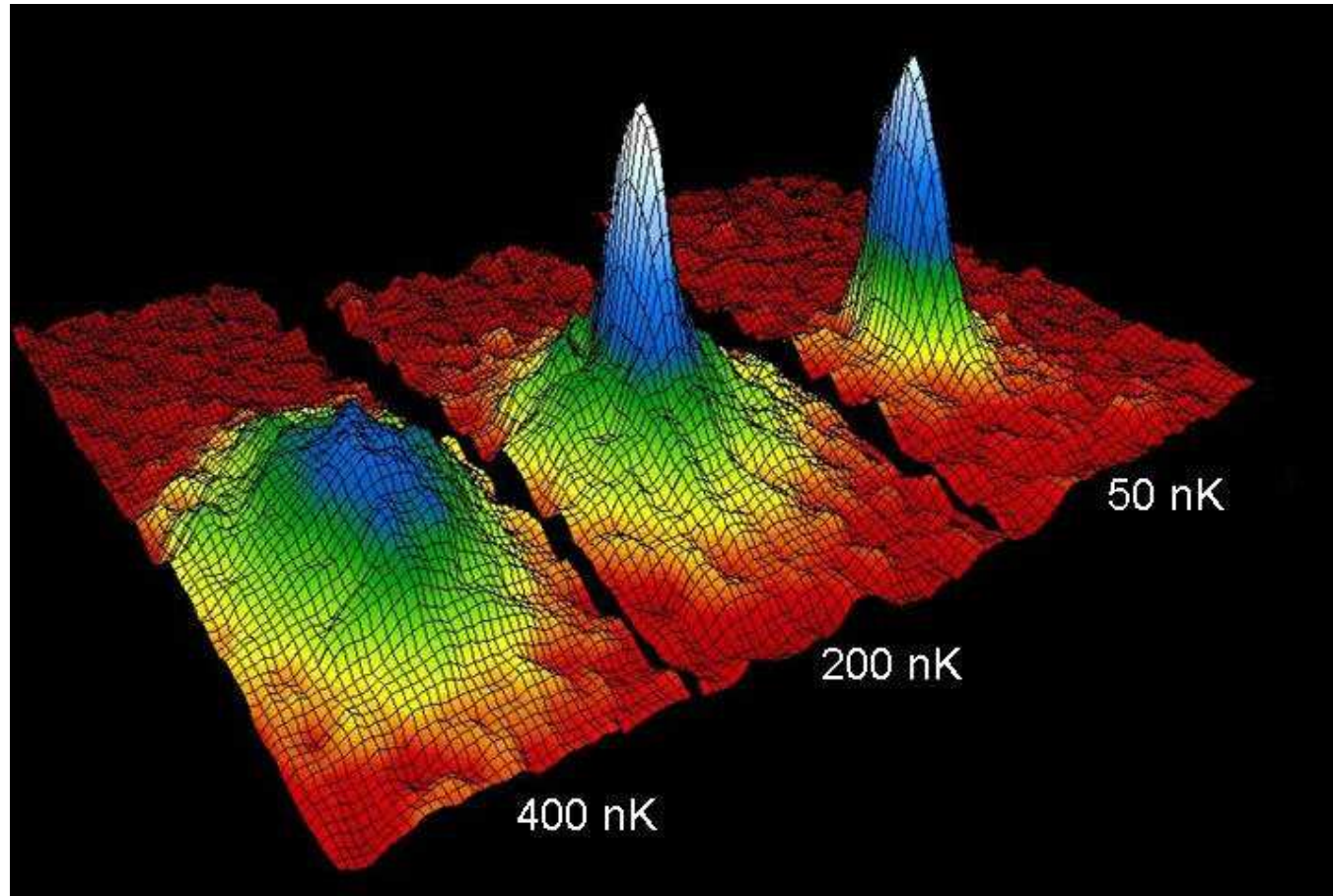
Robert Graham and Axel Pelster



1. Introduction: Ultracold Quantum Gases
2. Experimental Realizations of Dirty Bosons
3. Theoretical Description of Dirty Bosons
4. Huang-Meng Theory ( $T=0$ )
5. Shift of Condensation Temperature
6. Hartree-Fock Mean-Field Theory
7. Summary and Outlook

**SFB/TR 12: Symmetries and Universality in Mesoscopic Systems**

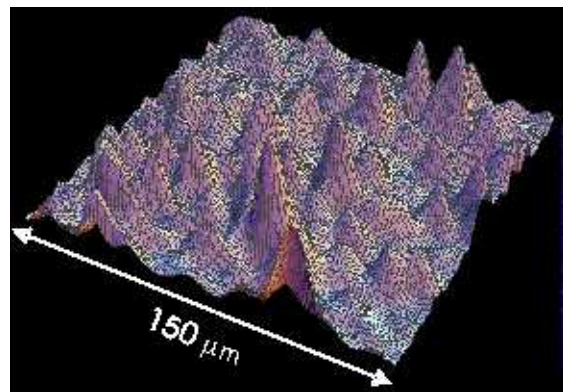
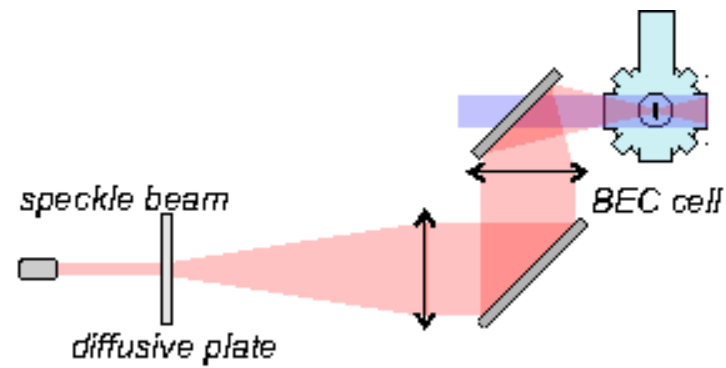
# 1 Introduction: Ultracold Quantum Gases



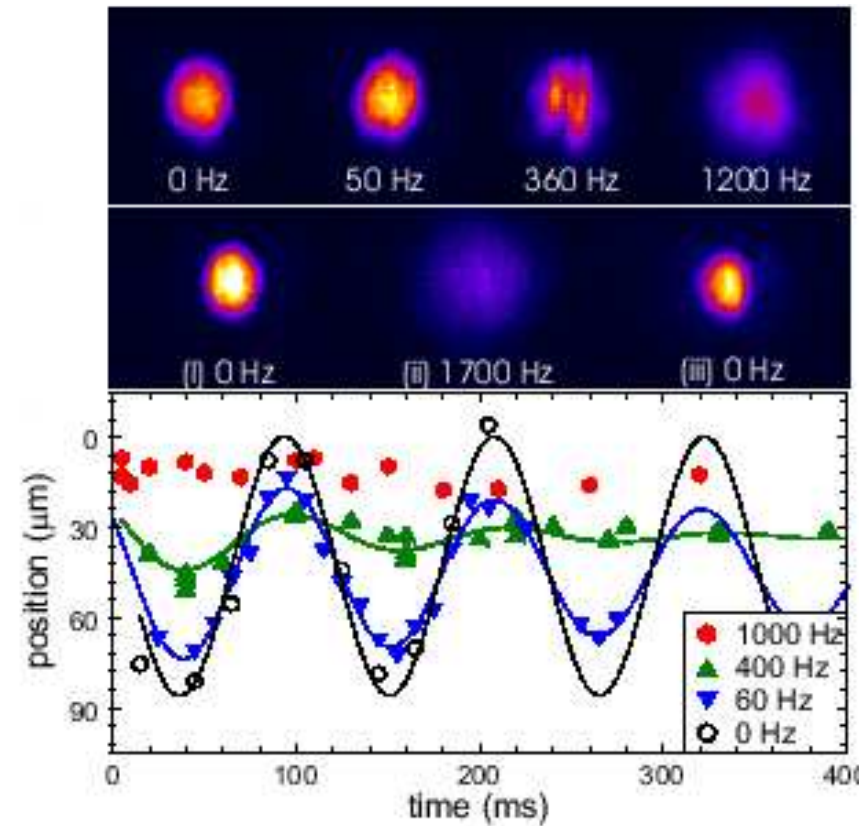
JILA (1995):  ${}^{87}_{37}\text{Rb}$ ,  $N=20\,000$ ,  $\omega_1 = \omega_2 = \omega_3/\sqrt{8} = 2\pi \times 120$  Hz

## 2.1 Magneto-Optical Trap

### Laser Speckles:

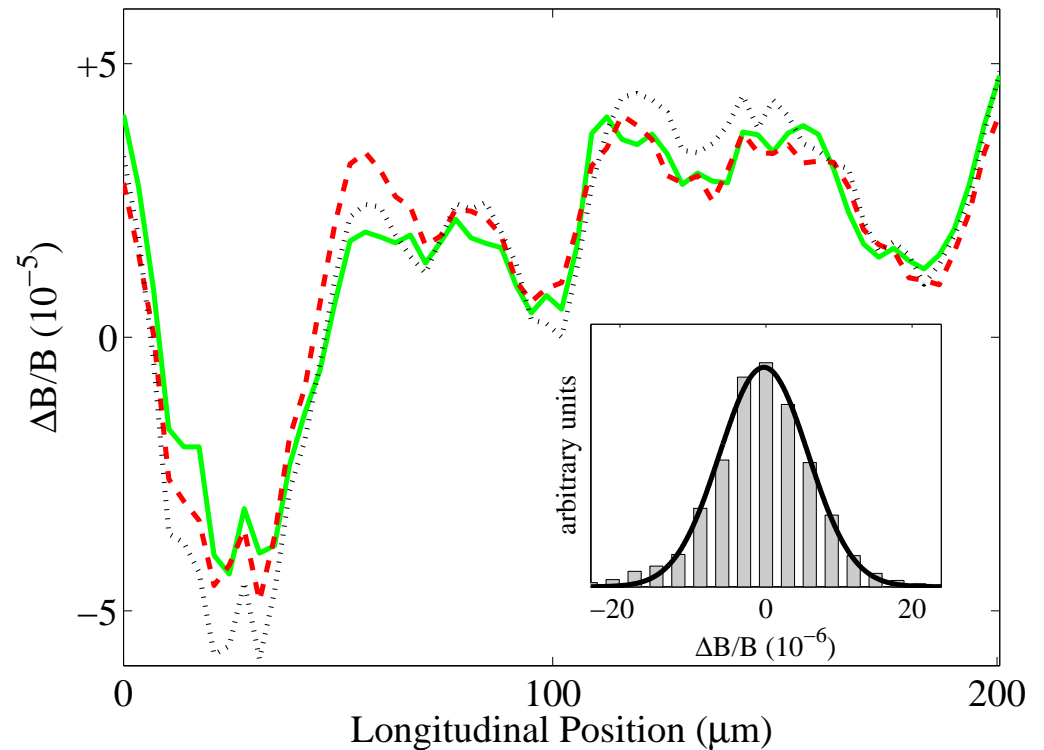
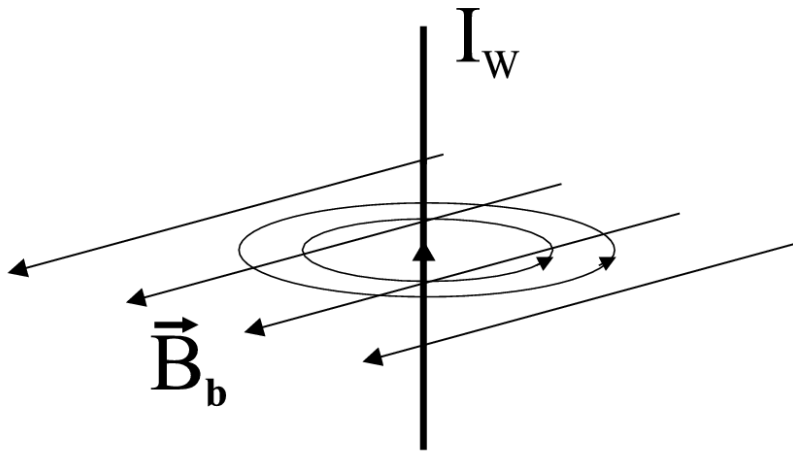


Inguscio *et al.*, PRL **95**, 070401 (2005)



global condensate vanishes

## 2.2 Wire Trap



**Distance:**  $d = 10 \mu\text{m}$

**Wire Width:**  $100 \mu\text{m}$

**Magnetic Field:** 10 G, 20 G, 30 G

**Deviation:**  $\Delta B/B \approx 10^{-4}$

Krüger *et al.*, Phys. Rev. A **76**, 063621 (2007)

Fortàgh and Zimmermann, Rev. Mod. Phys. **79**, 235 (2007)

## 3.1 Model System

### Action of a Bose Gas:

$$\mathcal{A} = \int_0^{\hbar\beta} d\tau \int d^3x \left\{ \psi^* \left[ \hbar \frac{\partial}{\partial \tau} - \frac{\hbar^2}{2M} \Delta + U(\mathbf{x}) + V(\mathbf{x}) - \mu \right] \psi + \frac{g}{2} \psi^{*2} \psi^2 \right\}$$

### Properties:

- harmonic trap potential:  $U(\mathbf{x}) = \frac{M}{2} \omega^2 \mathbf{x}^2$
- disorder potential:  $V(\mathbf{x})$ ; bounded from below, i.e.  $V(\mathbf{x}) \geq V_0$

$$\overline{V(\mathbf{x}_1)} = 0, \quad \overline{V(\mathbf{x}_1)V(\mathbf{x}_2)} = R(\mathbf{x}_1 - \mathbf{x}_2), \quad \dots$$

- chemical potential:  $\mu$
- repulsive interaction:  $g = \frac{4\pi\hbar^2 a}{M}$
- periodic Bose fields:  $\psi(\mathbf{x}, \tau + \hbar\beta) = \psi(\mathbf{x}, \tau)$

## 3.2 Grand-Canonical Potential

**Aim:**

$$\Omega = -\frac{1}{\beta} \overline{\ln \mathcal{Z}}$$
$$\mathcal{Z} = \oint D\psi D\psi^* e^{-\mathcal{A}[\psi^*, \psi]/\hbar}$$

**Problem:**

$$\overline{\ln \mathcal{Z}} \neq \ln \overline{\mathcal{Z}}$$

**Solution: Replica Trick**

$$\Omega = -\frac{1}{\beta} \lim_{N \rightarrow 0} \frac{\overline{\mathcal{Z}^N} - 1}{N}$$

## 3.3 Replica Trick

### Disorder Averaged Partition Function:

$$\overline{\mathcal{Z}^N} = \overline{\int \left\{ \prod_{\alpha=1}^N D\psi_{\alpha} D\psi_{\alpha}^* \right\} e^{-\sum_{\alpha=1}^N \mathcal{A}([\psi_{\alpha}^*, \psi_{\alpha}])/\hbar}} = \int \left\{ \prod_{\alpha=1}^N D\psi_{\alpha} D\psi_{\alpha}^* \right\} e^{-\mathcal{A}^{(N)}/\hbar}$$

### Replicated Action:

$$\begin{aligned} \mathcal{A}^{(N)} = & \int_0^{\hbar\beta} d\tau \int d^3x \sum_{\alpha=1}^N \left\{ \psi_{\alpha}^*(\mathbf{x}, \tau) \left[ \hbar \frac{\partial}{\partial \tau} - \frac{\hbar^2}{2M} \Delta + U(\mathbf{x}) - \mu \right] \psi_{\alpha}(\mathbf{x}, \tau) \right. \\ & \left. + \frac{g}{2} \psi_{\alpha}^*(\mathbf{x}, \tau)^2 \psi_{\alpha}(\mathbf{x}, \tau)^2 \right\} \\ & - \frac{1}{2\hbar} \int_0^{\hbar\beta} d\tau \int_0^{\hbar\beta} d\tau' \int d^3x \int d^3x' \sum_{\alpha=1}^N \sum_{\alpha'=1}^N \\ & \times \psi_{\alpha}^*(\mathbf{x}, \tau) \psi_{\alpha}(\mathbf{x}, \tau) R(\mathbf{x} - \mathbf{x}') \psi_{\alpha'}^*(\mathbf{x}', \tau') \psi_{\alpha'}(\mathbf{x}', \tau') + \dots \end{aligned}$$

**Similar:** disorder averaged correlation functions

## 4.1 Condensate Density

### Assumptions:

homogeneous Bose gas:  $U(\mathbf{x}) = 0$

$\delta$ -correlated disorder:  $R(\mathbf{x}) = R \delta(\mathbf{x})$

### Bogoliubov Theory:

background method:  $\psi_\alpha(\mathbf{x}, \tau) = \Psi_\alpha + \delta\psi_\alpha(\mathbf{x}, \tau)$

replica symmetry:  $\Psi_\alpha = \sqrt{n_0}$

### Result:

$$n_0 = n - \frac{8}{3\sqrt{\pi}} \sqrt{a n_0}^3 - \frac{M^2 R}{8\pi^{3/2} \hbar^4} \sqrt{\frac{n_0}{a}}$$

Huang and Meng, PRL **69**, 644 (1992)

Falco, Pelster, and Graham, Phys. Rev. A **75**, 063619 (2007)



## 4.2 Superfluid Density

### Galilei Boost:

$$\Delta\mathcal{A} = \int_0^{\hbar\beta} d\tau \int d^3x \psi^*(\mathbf{x}, \tau) \mathbf{u} \frac{\hbar}{i} \nabla \psi(\mathbf{x}, \tau)$$

$$d\Omega = -S dT - p dV - N d\mu - \mathbf{p} d\mathbf{u}$$

$$\mathbf{p} = - \left. \frac{\partial\Omega(T, V, \mu, \mathbf{u})}{\partial\mathbf{u}} \right|_{T, V, \mu} = MV n_n \mathbf{u} + \dots$$

### Result:

$$n_s = n - n_n = n - \frac{4}{3} \frac{M^2 R}{8\pi^{3/2} \hbar^4} \sqrt{\frac{n_0}{a}}$$

Huang and Meng, PRL **69**, 644 (1992)

Falco, Pelster, and Graham, Phys. Rev. A **75**, 063619 (2007)

## 4.3 Collective Excitations

### Hydrodynamic Equation in Trap With Disorder:

$$\begin{aligned} m \frac{\partial^2}{\partial t^2} \delta n(\mathbf{x}, t) - \nabla \left[ g n_{\text{TF}}(\mathbf{x}) \nabla \delta n(\mathbf{x}, t) \right] \\ = -\nabla^2 \left[ 3g n_R(\mathbf{x}) \delta n(\mathbf{x}, t) \right] - \nabla \left[ \frac{4g}{3} n_R(\mathbf{x}) \nabla \delta n(\mathbf{x}, t) \right] \end{aligned}$$

$n_R(\mathbf{x})$  : Huang-Meng depletion in trap

$n_{\text{TF}}(\mathbf{x}) = [\mu - V(\mathbf{x})] / g$  : Thomas-Fermi density

### Violation of Kohn Theorem:

Surface dipole mode  
( $n = 0, l = 1$ ):

$$\frac{\delta \omega_{\text{dip}}(\xi = 0)}{\omega_{\text{dip}}} = -\frac{5\pi}{16} \frac{M^2 R}{8\pi^{3/2} \hbar^4 \sqrt{n_{\text{TF}}(\mathbf{0})} a}$$

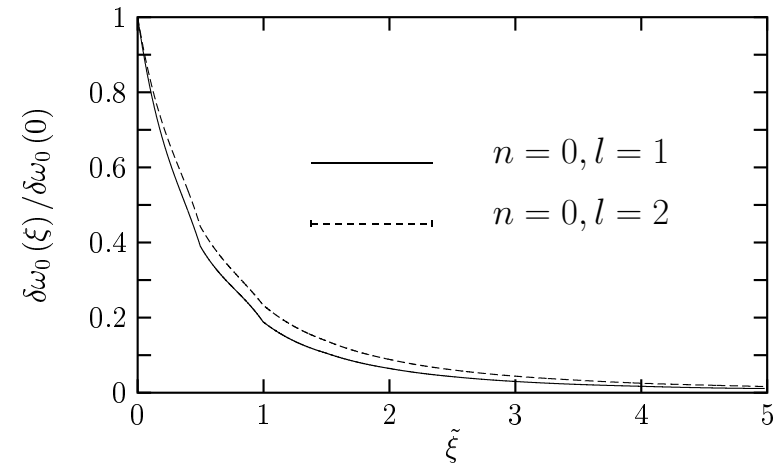
Falco, Pelster, and Graham, Phys. Rev. A **76**, 013624 (2007)

## 4.4 Comparison With Experiment

### Typical Values:

Inguscio *et al.*, PRL **95**, 070401 (2005)

$$\left. \begin{array}{l} \xi = 10 \mu\text{m} \\ R_{\text{TF}} = 100 \mu\text{m} \\ l_{\text{HO}} = 10 \mu\text{m} \end{array} \right\} \tilde{\xi} = \frac{\xi R_{\text{TF}}}{l_{\text{HO}}^2 \sqrt{2}} \approx 7$$



⇒ **Disorder effect vanishes in laser speckle experiment**

### Improvement:

laser speckle setup with correlation length  $\xi = 1 \mu\text{m}$

Aspect *et al.*, New J. Phys. **8**, 165 (2006)

⇒ **Disorder effect should be measurable**

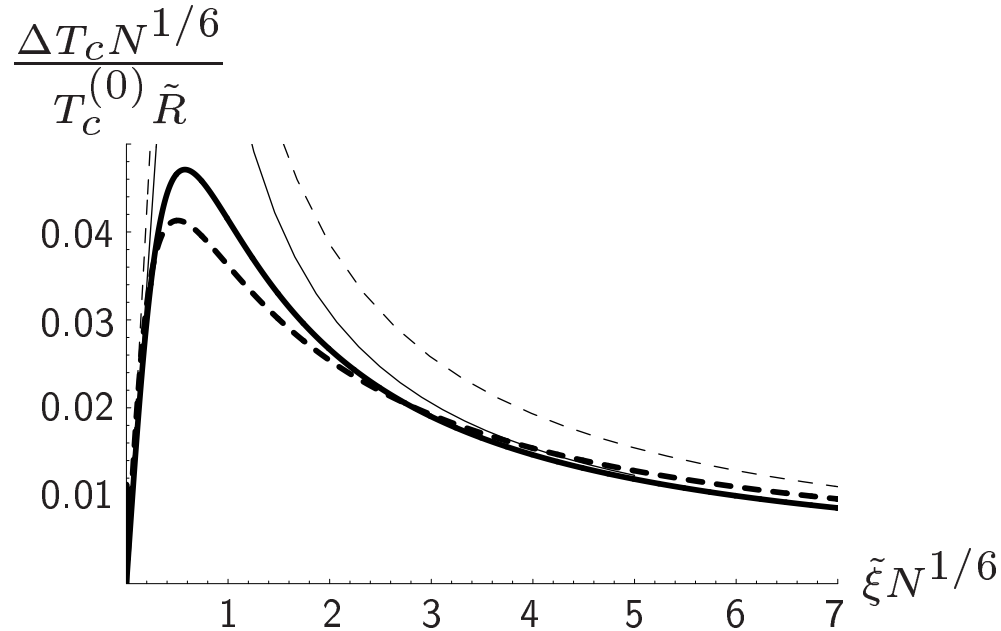
Falco, Pelster, and Graham, Phys. Rev. A **76**, 013624 (2007)

## 5.1 Earlier Results

trapped Bose gas	homogeneous Bose gas
$T_c^{(0)} = \frac{\hbar\omega_g}{k_B} \left[ \frac{N}{\zeta(3)} \right]^{1/3}$	$T_c^{(0)} = \frac{2\pi\hbar^2}{k_B M} \left[ \frac{n}{\zeta(3/2)} \right]^{2/3}$
$\frac{\Delta T_c}{T_c^{(0)}} = -3.426 \frac{a}{\lambda_c^{(0)}}$ <p>Giorgini et al., PRA <b>54</b>, R4633 (1996)            Gerbier et al., PRL <b>92</b>, 030405 (2004)</p>	$\frac{\Delta T_c}{T_c^{(0)}} = 1.3 a n^{1/3}$ <p>Kleinert, Mod. Phys. Lett. B <b>17</b>, 1011 (2003)            Kastening, Phys. Rev. A <b>69</b>, 043613 (2004)</p>
$R(\mathbf{x}) = ?$ $\frac{\Delta T_c}{T_c^{(0)}} = ?$	$R(\mathbf{x}) = R \delta(\mathbf{x})$ $\frac{\Delta T_c}{T_c^{(0)}} = -\frac{M^2 R}{3\pi [\zeta(3/2)]^{2/3} \hbar^2 n^{1/3}}$ <p>Lopatin and Vinokur, PRL <b>88</b>, 235503 (2002)</p>

**Procedure:**  $n = n(\mu), \quad \mu \nearrow \mu_c \Rightarrow T_c$

## 5.2 Our Results



solid: Gaussian

dashed: Lorentzian

**Length Scale:**

$$l_{\text{HO}} = \sqrt{\frac{\hbar}{M\omega_g}} \quad , \quad \omega_g = (\omega_1\omega_2\omega_3)^{1/3}$$

**Dimensionless Units:**

$$\tilde{\xi} = \frac{\xi}{l_{\text{HO}}} \quad , \quad \tilde{R} = \frac{R}{\left(\frac{\hbar^2}{Ml_{\text{HO}}^2}\right)^2 l_{\text{HO}}^3} = \frac{M^{3/2}R}{\hbar^{7/2}\omega_g^{1/2}}$$

Timmer, Pelster, and Graham, Europhys. Lett. **76**, 760 (2006)

Klünder, Pelster, and Graham, to be published

## 6.1 Order Parameters

### Definition:

$$\lim_{|\mathbf{x}-\mathbf{x}'|\rightarrow\infty} \overline{\langle \psi(\mathbf{x}, \tau) \psi^*(\mathbf{x}', \tau) \rangle} = n_0$$
$$\lim_{|\mathbf{x}-\mathbf{x}'|\rightarrow\infty} \overline{|\langle \psi(\mathbf{x}, \tau) \psi^*(\mathbf{x}', \tau) \rangle|^2} = (n_0 + q)^2$$

### Note:

$q$  is similar to Edwards-Anderson order parameter of spin-glass theory

### Hartree-Fock Mean-Field Theory:

Self-consistent determination of  $n_0$  and  $q$  for  $R(\mathbf{x} - \mathbf{x}') = R \delta(\mathbf{x} - \mathbf{x}')$

### Phase Classification:

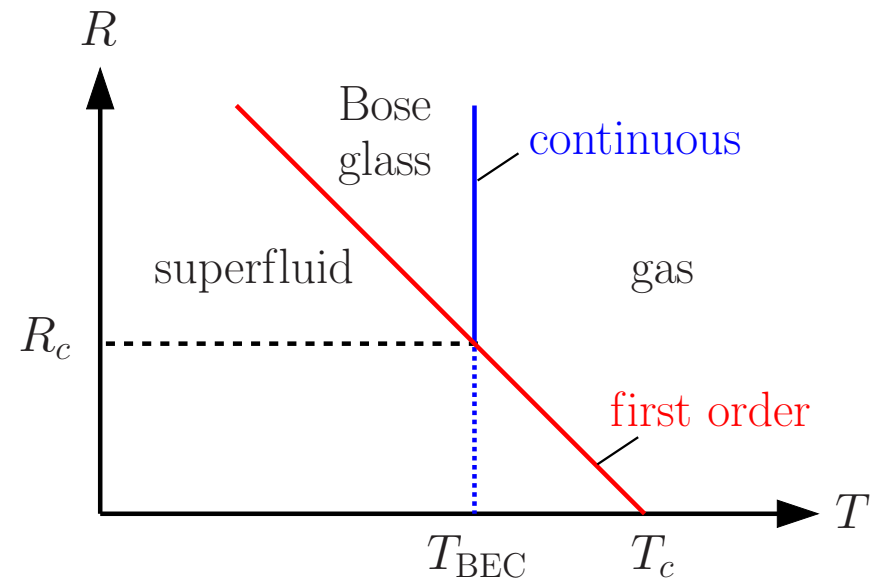
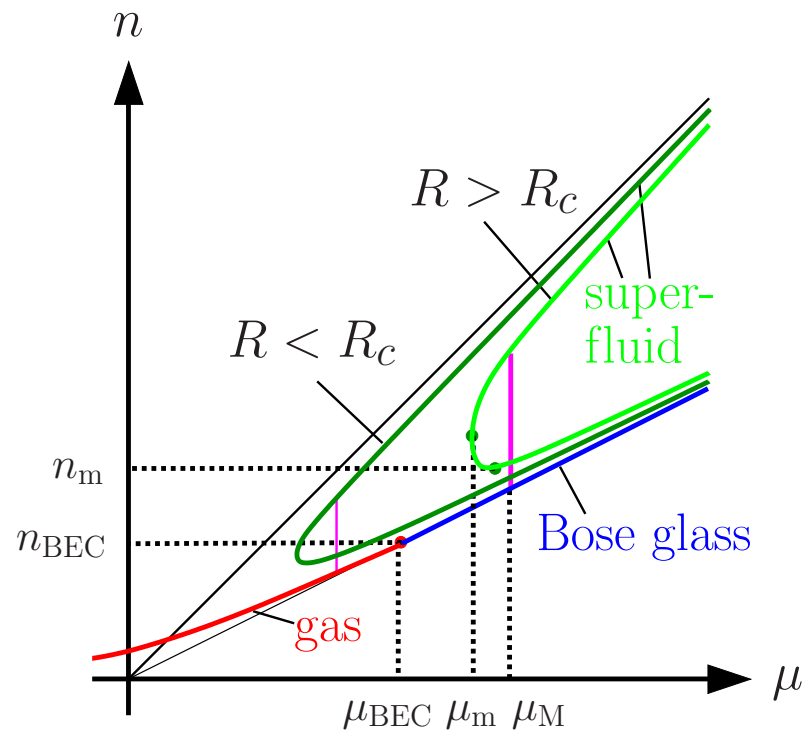
gas	Bose glass	superfluid
$q = n_0 = 0$	$q > 0, n_0 = 0$	$q > 0, n_0 > 0$

## 6.2 Hartree-Fock Results

**Isotherm:**  $T = \text{const.}$

**Phase Diagram:**  $\mu = \text{const.}$

disorder strength  $R = \text{const.}$



Graham and Pelster, Int. J. Bif. Chaos (in press)

## 7 Summary and Outlook

- **Frozen Disorder Potential:**

arises both artificially (laser speckles) or naturally (wire trap)

- **Bosons:**

local condensates in minima + global condensate + thermally excited

- **Localization Versus Transport:**

disorder reduces superfluidity

- **Phase Diagram:**

yet unknown for strong disorder

Navez, Pelster, and Graham, Appl. Phys. B **86**, 395 (2007)

- **Disordered Bosons in Lattice:**

Bose Glass versus Mott phase

Krutitsky, Pelster, and Graham, New J. Phys. **8**, 187 (2006)



## 8 Acknowledgements

<b>University Duisburg-Essen AG Robert Graham</b>	<b>Free University of Berlin AG Hagen Kleinert</b>
Giovanni Falco (Köln) Ben Klünder (LMU) Konstantin Krutitsky Patrick Navez (Leuven) Wieland Ronalter Matthias Timmer	Konstantin Glaum (Ulm) Alexander Hoffmann (LMU) Sebastian Kling (Bonn) Aristeu Lima (DAAD) Pascal Mattern Matthias Ohliger (Potsdam) Steffen Röthel (Münster) Ednilson Santos (DAAD) Sebastian Schmidt (Yale) Moritz Schütte (Golm) Parvis Soltan-Panahi (Hamburg)
<b>SFB/TR 12: Symmetries and Universality in Mesoscopic Systems</b>	<b>SPP 1116: Interactions in Ultra-Cold Atomic and Molecular Gases</b>

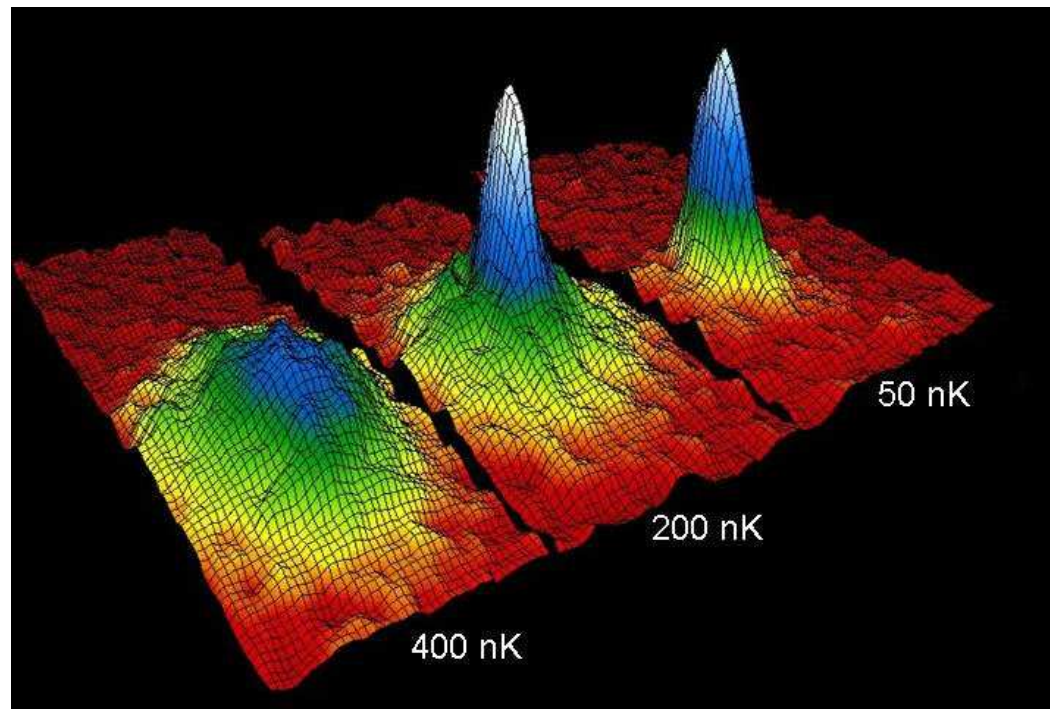
## 422nd Wilhelm and Else Heraeus Seminar

# Quo Vadis BEC?

Bad Honnef, October 29 - 31, 2008

*Scientific Coordinators:* Martin Holthaus (Oldenburg) and Axel Pelster (Duisburg-Essen)

*Research Topics:* BEC/BCS Crossover, Dipolar Gases, **Disorder**, Dynamics, Quantum Information, Spinor Bose and Fermi Gases, Strong Correlations, Tunneling



[http://www.theo-phys.uni-essen.de/tp/ags/pelster\\_dir/Heraeus/index.html](http://www.theo-phys.uni-essen.de/tp/ags/pelster_dir/Heraeus/index.html)