

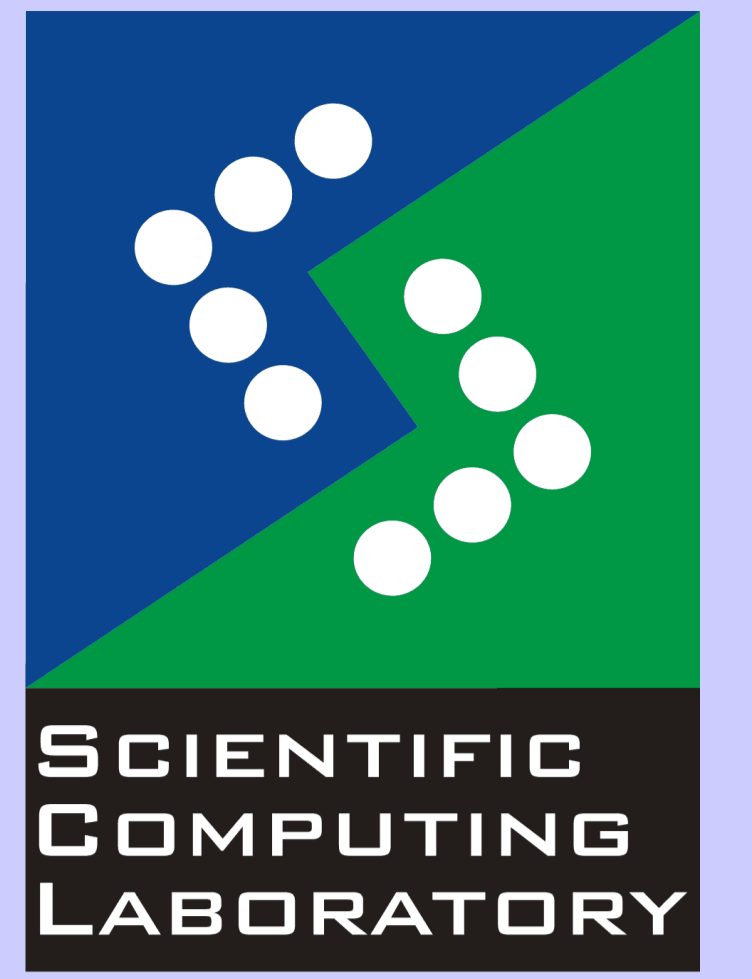
Influence of interface roughness on relaxation rates and optical gain in a quantum cascade laser

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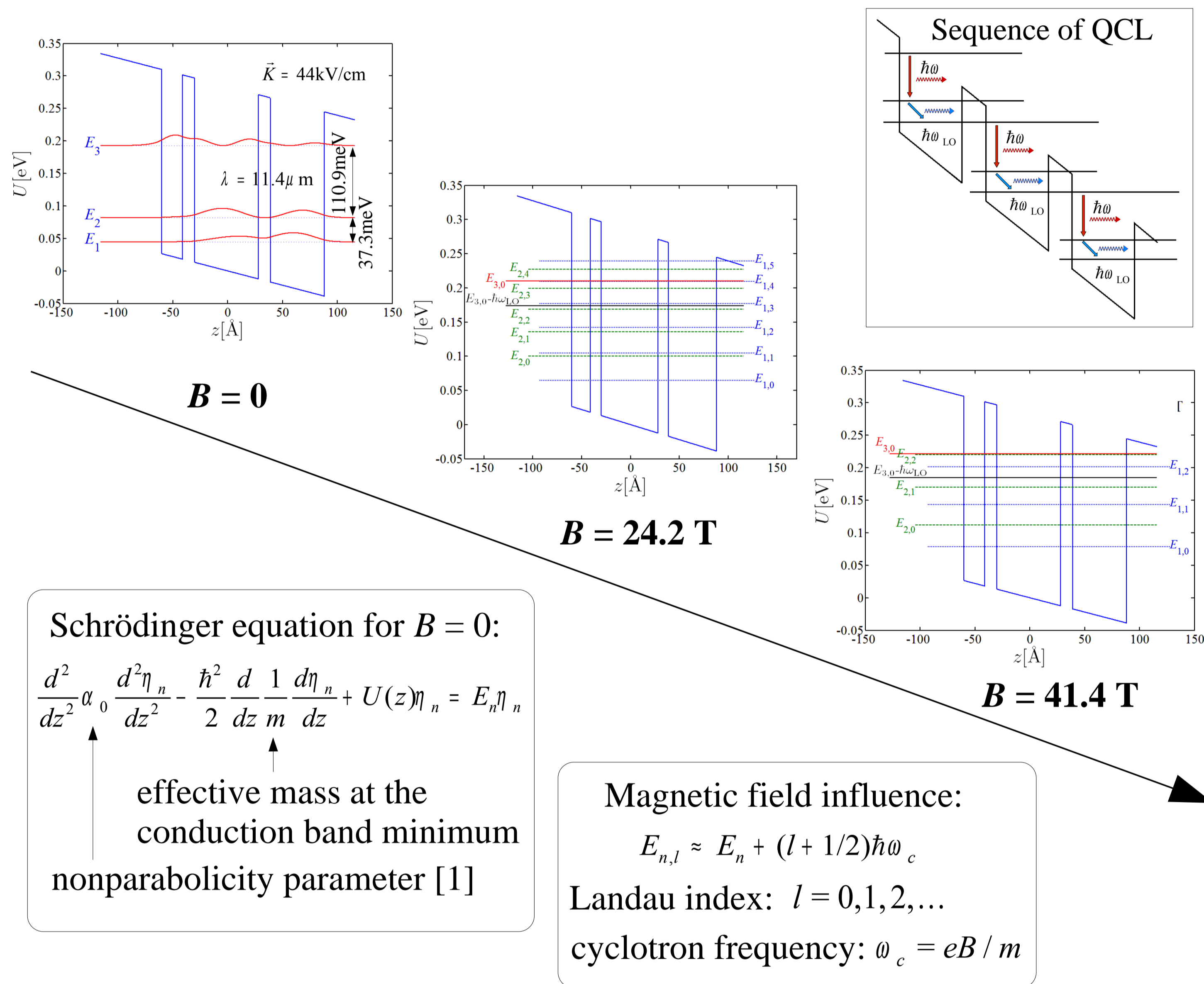
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Introduction: We present a detailed theoretical analysis of LO-phonon and interface roughness scattering influence on the operation of GaAs/AlGaAs quantum cascade laser (QCL) in the presence of an intense external magnetic field. We observe strong variations in the life time of the upper state, population inversion and optical gain when magnetic field is increased. The positions and magnitude of peaks are found to be a result of the combined action of both studied scattering mechanisms and strongly influenced by temperature. Still, incorporation of the interface roughness scattering mechanism into the model did not create new resonant peaks of the optical gain. We show that the decrease in the optical gain at elevated temperatures, is moderated by the occurrence of interface roughness scattering, which remains unchanged with increasing temperature.

(1) Quantum cascade laser (QCL) in magnetic field



(2) Electron scattering rates

1) electron - LO - phonon scattering emission, absorption [1]:

$$\frac{1}{\tau_{(n_i, l_i) \rightarrow (n_f, l_f)}^{LO}} = \frac{2\pi}{\hbar} \sum_{\vec{q}} \left| \langle n_f, l_f, k_{x_f}, n_q \pm 1 | \hat{H}_{e-ph}(\vec{q}) | n_i, l_i, k_{x_i}, n_q \rangle \right|^2 J^{LO}$$

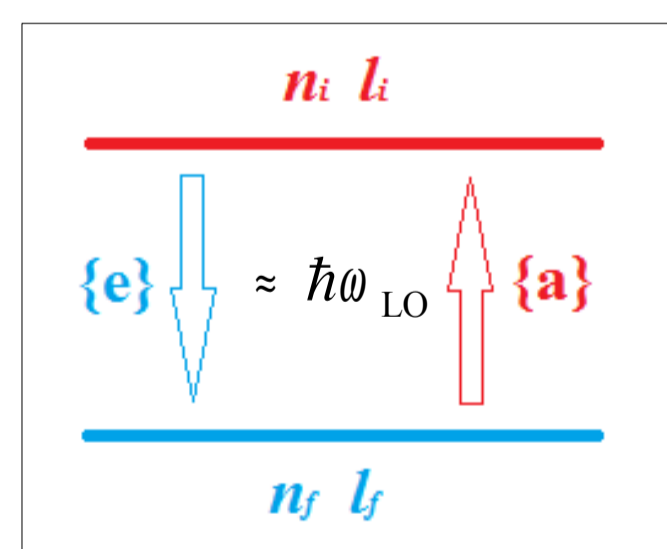
$$\hat{H}_{e-ph}(\vec{q}) = \sum_{\vec{q}} i \frac{g}{q} \left(e^{-i\vec{q}\cdot\vec{r}} a_{\vec{q}}^+ - e^{i\vec{q}\cdot\vec{r}} a_{\vec{q}} \right)$$

phonon wave vector \vec{q} $g^2 \sim \hbar \omega_{LO}$ Frölich factor:

creation and annihilation operators

emission and absorption (temperature dependence):

$$\frac{1}{\tau_{(n_f, l_f) \rightarrow (n_i, l_i)}^{LO, \{a\}}} = \frac{1}{\tau_{(n_i, l_i) \rightarrow (n_f, l_f)}^{LO, \{e\}}} \frac{1}{e^{kT}}$$



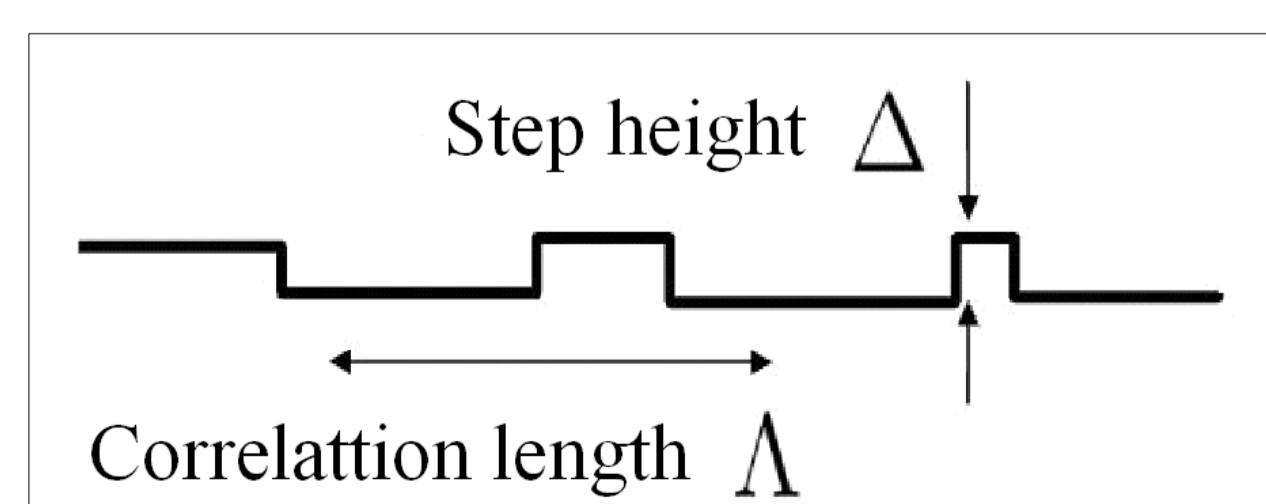
2) electron - interface roughness (IR) scattering [2]:

$$\left\langle \frac{1}{\tau_{(n_i, l_i) \rightarrow (n_f, l_f)}^{IR}}(z_i) \right\rangle = \frac{2\pi}{\hbar} \left\langle \sum_{k_{x_i}, k_{x_f}} \left| \langle n_f, l_f, k_{x_f} | \hat{H}_{IR} | n_i, l_i, k_{x_i} \rangle \right|^2 \right\rangle J^{IR}$$

$$\hat{H}_{IR} = U_0 \delta(z - z_i) \Delta(\vec{r})$$

Gaussian correlation function:

$$\langle \Delta(\vec{r}) \Delta(\vec{r}') \rangle = \Delta^2 e^{-\frac{|\vec{r}-\vec{r}'|^2}{\Lambda^2}}$$



3) total scattering (LO - phonon + interface roughness):

$$\frac{1}{\tau_{(n_i, l_i) \rightarrow (n_f, l_f)}} = \frac{1}{\tau_{(n_i, l_i) \rightarrow (n_f, l_f)}^{LO}} + \sum_{z_i} \left\langle \frac{1}{\tau_{(n_i, l_i) \rightarrow (n_f, l_f)}^{IR}}(z_i) \right\rangle$$

(3) Optical gain and rate equations

$$g_{3 \rightarrow 2} = \frac{2e^2 \pi^2}{\bar{n} \epsilon_0} \frac{d_{3 \rightarrow 2}^2}{\lambda} \sum_i \delta(E_{3,i} - E_{2,i} - \hbar \omega) (N_{3,i} - N_{2,i})$$

material refractive index \bar{n}
wavelength and the frequency of the emitted light λ
population inversion $(N_{3,i} - N_{2,i})$
transition matrix element: $d_{3 \rightarrow 2} = \int \eta_3(z) z \eta_2(z) dz$

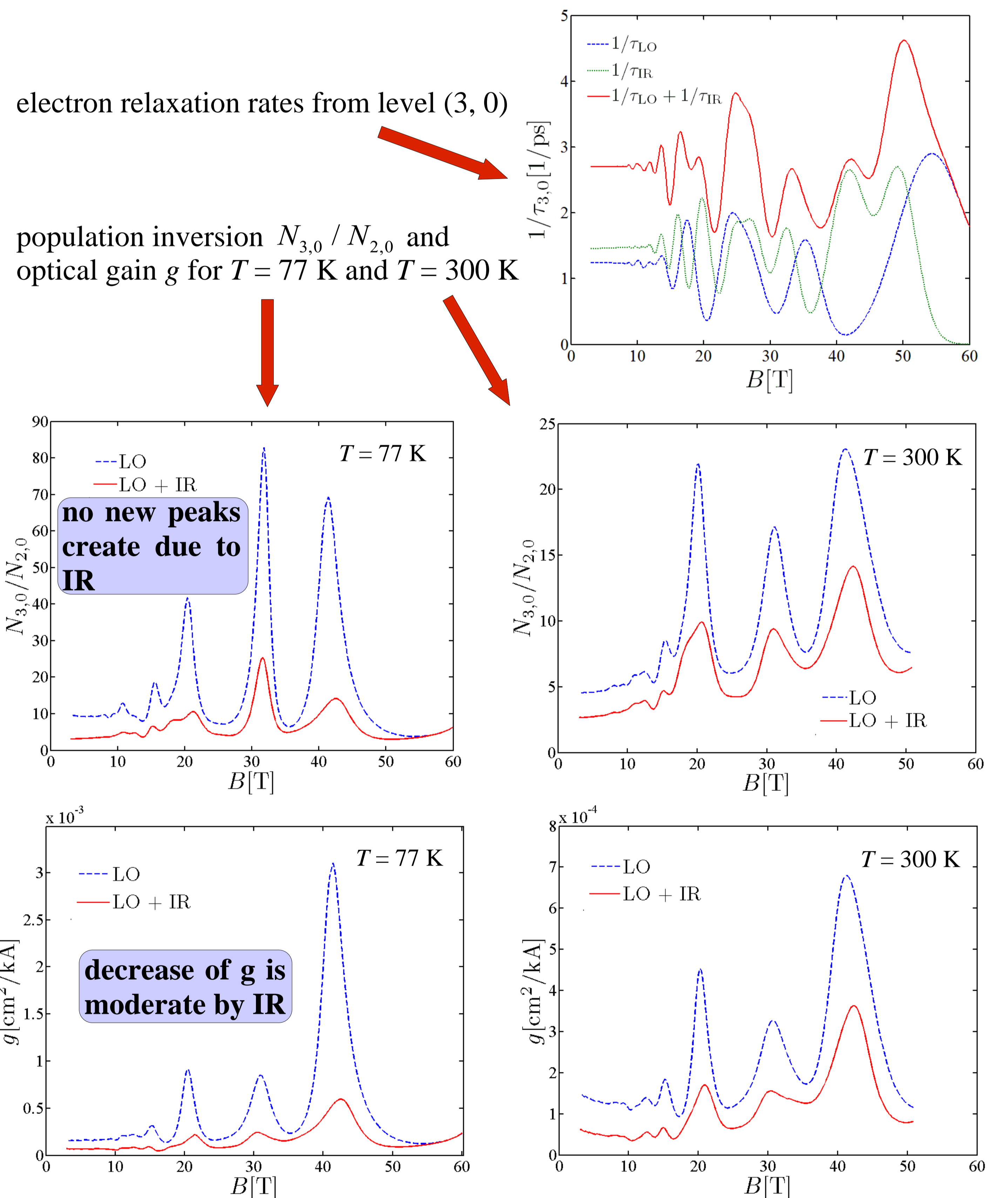
(4) Rate equations

extraction current J_i

$$N_i \sum_{j \neq i} \frac{\bar{f}_j}{\tau_{i \rightarrow j}} - \bar{f}_i \sum_{j \neq i} \frac{N_j}{\tau_{j \rightarrow i}} + \frac{J_i}{e} = 0, \quad \bar{f}_i = 1 - \frac{\pi \hbar}{eB} N_i$$

electron areal density N_i
probability that the state is not occupied according to the Fermi-Dirac distribution \bar{f}_i

(5) Results



References:

- [1] U. Ekenberg, *Phys. Rev. B* 40, 7714 (1989).
- [2] C. Becker, A. Vasanelli, C. Sirtori and G. Bastard, *Phys. Rev. B* 69, 115328 (2004).
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